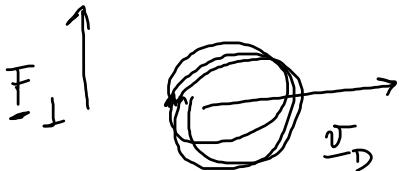


$$\underline{F} = \overrightarrow{\text{const}}$$

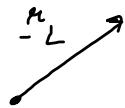
$\Rightarrow \underline{a}_{\parallel} = \underline{F}_{\parallel}/m \Rightarrow \underline{v}_{\parallel}(t) = \underline{v}_{0\parallel} + \underline{a}_{\parallel} t$

$\textcircled{\perp}$ Larmor motion + $\underline{v}_D = \frac{\underline{F}_{\perp} \times \underline{B}}{qB^2}$



Non uniform $\underline{E}(x, t)$ $\underline{B}(x, t)$

Sufficiently small variation



T_L - time scale

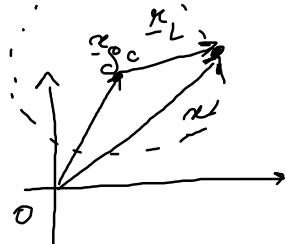
scale where $\Delta x \gg r_L$.

$\Delta t \gg T_L$

$|E|$ and $|B|$ change significantly

$$\frac{m \underline{\dot{x}}}{dt} = q \underline{\underline{E}}(\underline{x}, t) + q \underline{\underline{J}} \times \underline{\underline{B}}(\underline{x}, t)$$

Small variations:



Gyro-centre \underline{x}_{gc}

$$v_L \text{ Larmor radius}$$

$$\underline{v} = \underline{v}_{gc} + \underline{v}_L$$

$$\underline{x} = \underline{x}_{gc} + \underline{x}_L$$

$$\begin{aligned} & \underline{x}_L x \partial_x + \underline{x}_L y \partial_y \\ & + \underline{x}_L z \partial_z \end{aligned}$$

$$\underline{\underline{E}}(\underline{x}, t) = \underline{\underline{E}}\left(\underline{x}_{gc} + \underline{x}_L, t\right) \approx \underline{\underline{E}}\left(\underline{x}_{gc}\right) + (\underline{x}_L \cdot \nabla) \underline{\underline{E}}(\underline{x})$$

$$\underline{\underline{B}}(\underline{x}, t) \approx \underline{\underline{B}}(\underline{x}_{gc}) + (\underline{x}_L \cdot \nabla) \underline{\underline{B}} \Big|_{\underline{x} = \underline{x}_{gc}}$$

Taylor expansion

around \underline{x}_{gc} 1st order

$$(1) \quad m \frac{d\vec{v}_{gc}}{dt} + m \frac{d\vec{v}_L}{dt} = q \left[E(\vec{x}_{gc}) + (\vec{n}_L \cdot \vec{\nabla}) E(\vec{x}) \Big|_{\vec{x}_{gc}} + \left(\frac{v}{\gamma_{gc}} + \vec{v}_L \right) \times \right. \\ \left. \times \left(B(\vec{x}_{gc}) + (\vec{n}_L \cdot \vec{\nabla}) B(\vec{x}) \Big|_{\vec{x}_{gc}} \right) \right]$$

Larmor motion:

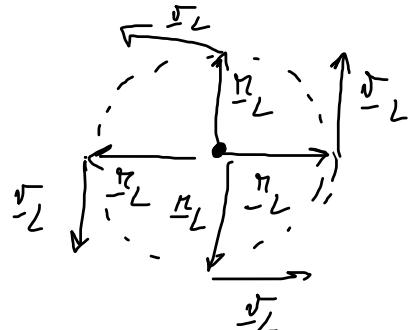
$$(2) \quad m \frac{d\vec{v}_L}{dt} = q \left(\vec{v}_L \times \underbrace{B(\vec{x}_{gc})} \right)$$

$$(1) - (2): \quad \langle m \frac{d\vec{v}_{gc}}{dt} \rangle = \langle q \left[E(\vec{x}_{gc}) + (\vec{n}_L \cdot \vec{\nabla}) E(\vec{x}) \Big|_{\vec{x}_{gc}} \right. \\ \left. + \vec{v}_{gc} \times B(\vec{x}_{gc}) + \vec{v}_{gc} \times \cancel{(\vec{n}_L \cdot \vec{\nabla}) B(\vec{x})} \Big|_{\vec{x}_{gc}} \right. \\ \left. - \vec{v}_L \times (\vec{n}_L \cdot \vec{\nabla}) B \Big|_{\vec{x}_{gc}} \right] \rangle$$

Time average over
a T_L : $\langle \rangle$

Time average $\leftarrow \rightarrow$

Linear terms $\left\{ \begin{array}{l} \underline{n}_L \\ \underline{\sigma}_L \end{array} \right. = 0$



After averaging:

$$m \frac{d\underline{\sigma}_{gc}}{dt} = q \bar{\underline{\sigma}}_{gc} + q \underline{\sigma}_{gc} \times \underline{B}(\underline{x}_{gc}) \\ + \langle q \underline{\sigma}_L \times (\underline{n}_L \cdot \underline{\sigma}) \underline{B} \rangle$$

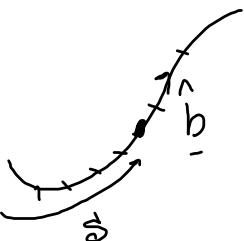
Separate // and \perp terms

$$\underline{\sigma}_{gc} = \underline{\sigma}_{\perp gc} + \underline{\sigma}_{\parallel gc}^{\wedge}$$

\hat{b} : vector associate to \underline{B}

$$\begin{aligned} \frac{d \underline{\underline{\sigma}}_{gc}}{dt} &= \frac{\partial}{\partial t} \left(\underline{\underline{\sigma}}_{gc\perp} \right) + \cancel{\frac{\partial}{\partial t} \left(\underline{\underline{\sigma}}_{gc} \hat{\underline{b}} \right)} \\ &= \frac{\partial \underline{\underline{\sigma}}_{gc\perp}}{\partial t} + \frac{\partial \underline{\underline{\sigma}}_{\parallel}}{\partial t} \hat{\underline{b}} + \cancel{\frac{\partial \hat{\underline{b}}}{\partial t}} \end{aligned}$$

$\frac{\partial}{\partial S} = (\hat{\underline{b}} \cdot \nabla)$



$$\hat{\underline{b}} = \hat{\underline{b}}(S)$$

$$\frac{\partial \hat{\underline{b}}}{\partial t} = \frac{\partial \hat{\underline{b}}}{\partial S} \cdot \underbrace{\frac{\partial S}{\partial t}}_{\nabla} = \nabla \cdot \frac{\partial \hat{\underline{b}}}{\partial S} = \nabla \cdot (\hat{\underline{b}} \cdot \nabla) \hat{\underline{b}}$$

$$\frac{d \underline{\underline{\sigma}}_{gc}}{dt} = \cancel{\frac{d \underline{\underline{\sigma}}_{gc\perp}}{dt}} + \frac{\partial \underline{\underline{\sigma}}_{\parallel}}{\partial t} \hat{\underline{b}} + \cancel{\frac{\partial \hat{\underline{b}}}{\partial t}} \hat{\underline{b}}$$

\parallel

$$m \frac{d\vec{v}_{\parallel}}{dt} = q \vec{E}_{\parallel} + q \vec{v}_{\perp} \times (\underline{\mu}_L \cdot \underline{\nabla}) \underline{B}_{\parallel}$$


\perp

$$m \frac{d\vec{v}_{\perp g_C}}{dt} = q \vec{J}_{g_C} \times \underline{B}(\underline{x}_{g_C}) + q \vec{E}_{\perp} - q \vec{v}_{\perp} \times (\underline{\mu}_L \cdot \underline{\nabla}) \underline{B}_{\perp}$$

$\underbrace{m \vec{v}_{\parallel} (\underline{b} \cdot \underline{\nabla}) \underline{b}}$

$$\frac{m d\vec{v}_{\perp g_C}}{dt} = q \vec{J}_{g_C} \times \underline{B}(\underline{x}_{g_C}) + \vec{F}_{\perp}$$

motion in a "uniform" magnetic field + a force

If the force is const:

$$\underline{v} = \underline{v}_L + \underline{v}_{g_C} \quad \underline{v}_{g_C} \neq \frac{\vec{F}_{\perp} \times \underline{B}}{q B^2}$$

$$\frac{\underline{v}}{D} = \underbrace{\frac{\underline{v}^{(0)}}{D}}_{\begin{array}{l} \cancel{\underline{F}_1 \times \underline{B}} \\ \cancel{q \underline{B}^2} \\ \cdot \end{array}} + \frac{\underline{v}^{(1)}}{D}$$

$\left| \frac{\underline{v}^{(1)}}{D} \right| < \left| \frac{\underline{v}^{(0)}}{D} \right|$

$$m \frac{d}{dt} \left(\frac{\underline{v}^{(0)}}{D} + \frac{\underline{v}^{(1)}}{D} \right) = q \left[\left(\frac{\underline{v}^{(0)}}{D} + \frac{\underline{v}^{(1)}}{D} \right) \times \underline{B} \right] + \underline{F}_1$$

$$\frac{\frac{d \underline{v}^{(1)}}{D}}{dt} \ll \frac{\frac{d \underline{v}^{(0)}}{D}}{dt}$$

$$\times \underline{B}$$

$$m \frac{d \underline{v}^{(0)}}{dt} \approx q \underbrace{\frac{\underline{v}^{(0)}}{D} \times \underline{B}}_{\cancel{\underline{F}_1 \times \underline{B}}} + q \underbrace{\frac{\underline{v}^{(1)}}{D} \times \underline{B}}_{\cancel{\underline{F}_1}} + \cancel{\underline{F}_1} \times \underline{B}$$

$$(\underline{C} \times \underline{B}) \times \underline{A} = (\underline{A} \cdot \underline{C}) \underline{B} - (\underline{A} \cdot \underline{B}) \underline{C} \quad \left| \begin{array}{l} (\underline{F}_1 \times \underline{B}) \times \underline{B} \\ \cancel{(\underline{F}_1 \times \underline{B}) \times \underline{B}} - \underline{B}^2 \underline{F}_1 \end{array} \right. = - \underline{B}^2 \underline{F}_1 = - \frac{\underline{F}_1}{q}$$

$$\begin{aligned}
 m \frac{d \underline{v}_D^{(0)}}{dt} \times \underline{B} &= q (\underline{r}_D^{(1)} \times \underline{B}) \times \underline{B} \\
 &= q (\underline{r}_D^{(1)} \cdot \underline{B}) \underline{B} - B^2 \underline{r}_D^{(1)} q \\
 &\equiv
 \end{aligned}$$

$$\boxed{\underline{v}_D^{(1)} = - \frac{m}{qB^2} \frac{d \underline{v}_D^{(0)}}{dt} \times \underline{B}}$$

$$\underline{r}_D^{(0)} = \frac{\underline{F}_\perp \times \underline{B}}{qB^2}$$

$$\begin{aligned}
 \underline{v}_D^{(1)} &= - \frac{m}{qB^2} \frac{d}{dt} \left[\frac{[\underline{F}_\perp \times \underline{B}]}{qB^2} \right] \times \underline{B} \longrightarrow - B^2 \frac{d \underline{F}_\perp}{dt} \\
 \underline{r}_D^{(0)} &= \frac{m}{qB^2} \frac{d \underline{F}_\perp}{dt}
 \end{aligned}$$

$$\oint \vec{F}_J = q \vec{E}_J$$

$$\vec{J}_D^{(1)} = \frac{m}{qB^2} \frac{d\vec{E}_J}{dt}$$

$$\vec{v}_D^{(1)} \parallel \vec{E}_J$$

$\vec{J}_D^{(1)}$ is charge dependent

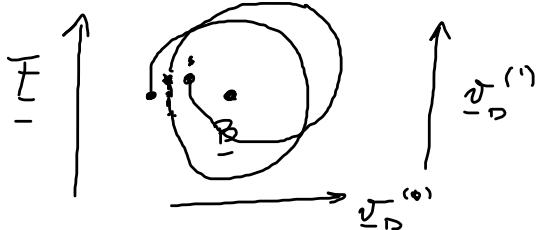
\Rightarrow it drives a current
polarization drift

$$\vec{E} = \text{const}$$

$$\vec{E} \uparrow$$



\vec{E} increases as a function of time



$$q \left(\underline{\omega}_L \times (\underline{n}_L \cdot \nabla) \underline{B} \right) = ?$$

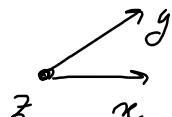
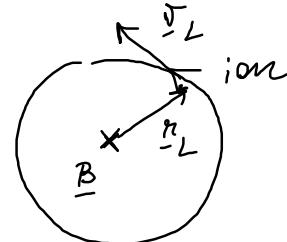
$$\underline{B} = B_x \hat{i} + B_y \hat{j} + B_z \hat{k}$$

$$\underline{\omega}_L = \omega_L \begin{bmatrix} -\sin(\omega_L t) \hat{i} \\ +\cos(\omega_L t) \hat{j} \end{bmatrix}$$

$$\underline{n}_L = \frac{\underline{\omega}_L}{|\underline{\omega}_L|} \begin{bmatrix} \cos(\omega_L t) \hat{i} \\ +\sin(\omega_L t) \hat{j} \end{bmatrix}$$

$$\frac{\omega_L}{\omega_L} \rightarrow \frac{\omega_L}{\omega_L} = \frac{\omega_L m}{qB} = q_L$$

$$q_L q \underline{\omega}_L \left[-\sin(\omega_L t) \hat{i} + \cos(\omega_L t) \hat{j} \right] \times \left(\begin{bmatrix} \cos(\omega_L t) \hat{i} \\ +\sin(\omega_L t) \hat{j} \end{bmatrix} \cdot \frac{\partial \hat{i}}{\partial x} \hat{i} + \frac{\partial \hat{j}}{\partial y} \hat{j} \right. \\ \left. + \frac{\partial \hat{k}}{\partial z} \hat{k} \right) \underline{B} >$$



$$= q \eta_1 v_1 \left\langle \left(-\sin(\omega_L t) \hat{i} + \cos(\omega_L t) \hat{j} \right) \times \right.$$

$$\left. \left(\cos(\omega_L t) \left(\frac{\partial B_x}{\partial x} \hat{i} + \frac{\partial B_y}{\partial x} \hat{j} + \frac{\partial B_z}{\partial x} \hat{k} \right) + \sin(\omega_L t) \left(\frac{\partial B_x}{\partial y} \hat{i} + \frac{\partial B_y}{\partial y} \hat{j} + \frac{\partial B_z}{\partial y} \hat{k} \right) \right) \right\rangle$$