

$$L(q_j, \dot{q}_j, t)$$

$$j = 1 \dots N$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_j} - \frac{\partial L}{\partial q_j} = 0$$

$$j = 1 \dots N$$

Exact invariants : if  $L$  does not depend on  $q_i \Rightarrow \frac{\partial L}{\partial q_i} = 0$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} = 0 \Rightarrow p_i = \frac{\partial L}{\partial \dot{q}_i} = \text{const}$$

→ conjugate momentum

If there is rotational symmetry:  $\mathcal{L}$  cannot depend on  $\varphi$

$$\Rightarrow \frac{\partial \mathcal{L}}{\partial \dot{\varphi}} = \text{const}$$

Strategy: 1) Write down or guess form  $\mathcal{L}$

2) Show that  $\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{q}_j} - \frac{\partial \mathcal{L}}{\partial q_j} = 0$  is the Lorentz force

$$\mathcal{L} = \frac{m\dot{\underline{v}}^2}{2} + q \underline{v} \cdot \underline{A}(\underline{x}, t) - q\phi(\underline{x}, t)$$

$\underline{A}$ : vector potential

$q_j$ :  $x, y, z$

$$\underline{E} = -\underline{\nabla}\phi - \partial \underline{A} / \partial t$$

$\phi$ : scalar potential

$$\underline{B} = \underline{\nabla} \times \underline{A}$$

$$q_j: x, y, z$$

$$\dot{q}_j: \dot{x}, \dot{y}, \dot{z}$$

$$\frac{\partial \mathcal{L}}{\partial x_j} = \frac{\partial}{\partial x_j} \sum_{i=1}^3 q v_i A_i(x, t) - q \frac{\partial \phi}{\partial x_j} =$$

$$\equiv \sum_{i=1}^3 q v_i \frac{\partial A_i}{\partial x_j} - q \frac{\partial \phi}{\partial x_j} =$$

$$\frac{\partial \mathcal{L}}{\partial v_j} = \frac{m}{2} \frac{\partial}{\partial v_j} \sum_{i=1}^3 v_i^2 + q \frac{\partial}{\partial v_j} \sum_{i=1}^3 v_i A_i = \frac{m}{2} 2v_j + q A_j =$$

$$= m v_j + q A_j(x, t)$$

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial v_j} = m \dot{v}_j + q \sum_{i=1}^3 \frac{\partial A_i}{\partial x_i} \dot{x}_i + q \frac{\partial A_j}{\partial t}$$

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{v}_j} - \frac{\partial \mathcal{L}}{\partial x_j} = 0$$

$$m \dot{v}_j + q \sum_{i=1}^3 \frac{\partial A_j}{\partial x_i} \dot{x}_i + q \frac{\partial A_j}{\partial t} - q \sum_{i=1}^3 \dot{x}_i \frac{\partial A_i}{\partial x_j} + \frac{\partial \phi}{\partial x_j} q = 0$$

$$m \underline{\underline{a}} = q (\underline{\underline{E}} + \underline{\underline{v}} \times \underline{\underline{B}})$$

$$m \dot{v}_j = -q \frac{\partial \phi}{\partial x_j} - q \frac{\partial A_j}{\partial t} + q \sum_{i=1}^3 \left( \dot{x}_i \frac{\partial A_i}{\partial x_j} - \frac{\partial A_j}{\partial x_i} \dot{x}_i \right)$$

$$+ q \left[ \begin{array}{c} -\frac{\partial \phi}{\partial x_j} \\ -\frac{\partial A_j}{\partial t} \end{array} \right]$$

↗  $\underline{\underline{E}}_j$

$$\underline{\underline{E}} = -\underline{\underline{\nabla}} \phi - \frac{\partial \underline{\underline{A}}}{\partial t}$$

$$m \dot{v}_j = q E_j + q \underbrace{\sum_{i=1}^3 v_i \left( \frac{\partial A_i}{\partial x_j} - \frac{\partial A_j}{\partial x_i} \right)}$$

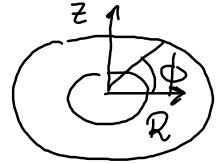
$$q (\underline{v} \times \underline{B})_j = q \left[ \underline{v} \times (\underline{\nabla} \times \underline{A}) \right]_j$$

$$\underline{v} \times (\underline{\nabla} \times \underline{A}) = \begin{pmatrix} v_x \\ v_y \\ v_z \end{pmatrix} \times \begin{bmatrix} \partial_x \\ \partial_y \\ \partial_z \end{bmatrix} \times \begin{bmatrix} A_x \\ A_y \\ A_z \end{bmatrix} = \begin{pmatrix} v_x \\ v_y \\ v_z \end{pmatrix} \times \begin{pmatrix} \partial_y A_z - \partial_z A_y \\ \partial_z A_x - \partial_x A_z \\ \partial_x A_y - \partial_y A_x \end{pmatrix}$$

$$\begin{pmatrix} v_y (\partial_x A_y - \partial_y A_x) - v_z (\partial_z A_x) \\ v_x (\partial_x A_y - \partial_y A_x) - v_z (\partial_z A_x) \\ -v_x A_z \end{pmatrix}$$

$$q \sum_{i=1}^3 v_i \left( \frac{\partial A_i}{\partial x_j} - \frac{\partial A_j}{\partial x_i} \right)$$

$$j=1$$



$$= q v_x \cancel{\frac{\partial A_x}{\partial x} - \frac{\partial A_x}{\partial x}} + q v_y \left( \frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right) + v_z q \left( \frac{\partial A_z}{\partial x} - \frac{\partial A_x}{\partial z} \right)$$

y) There is toroidal symm.  
 $\frac{\partial \mathcal{L}}{\partial \dot{\phi}} = \text{const}$

$$\mathcal{L} = \frac{m v^2}{2} + q \vec{v} \cdot \vec{A}(\vec{x}, t) - q \phi(\vec{x}, t)$$

$$\mathcal{L} = \frac{m}{2} (v_R^2 + v_z^2 + v_\phi^2) + q (v_R A_R + v_z A_z + v_\phi A_\phi) - q \phi$$

(R, \phi, z)

$$\mathcal{L} = \frac{m}{2} (\dot{v}_R^2 + \dot{v}_z^2 + \dot{v}_\varphi^2) + q (A_R \dot{v}_R + A_z \dot{v}_z + A_\varphi \dot{v}_\varphi) - q\phi$$

$$\frac{\partial \mathcal{L}}{\partial \dot{\phi}} = \frac{m R^2 \dot{\phi}}{2} + q R A_\varphi \Rightarrow p_\varphi = \underline{\underline{m R^2 \dot{\phi}}} + \underline{\underline{q R A_\varphi}}$$

toroidal  
component

$$v_\varphi = R \dot{\phi}$$

$$[p_\varphi] = [m R \cdot \dot{v}_\varphi] = \text{angular momentum}$$

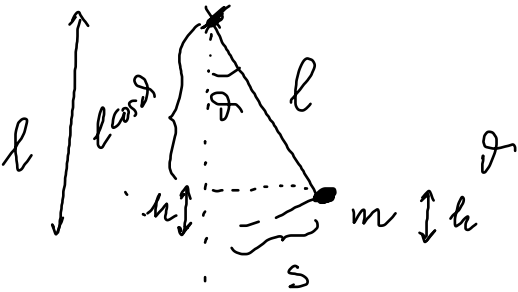
$$A_\varphi = R \dot{\phi}$$

$$v_\varphi = \frac{d(A_\varphi)}{dt} = R \dot{\phi}$$

Canonical

$$v_\varphi^2 = R^2 \dot{\phi}^2 \Rightarrow \frac{d}{d\phi} v_\varphi^2 = 2 \dot{\phi} R^2$$

# Simple pendulum



$\mathcal{L}$  = Kinetic en. - Potent. energy

$$K = \frac{1}{2} m v^2 = \frac{1}{2} m l^2 \dot{\vartheta}^2$$

$$v = \frac{ds}{dt} = \frac{d(l\vartheta)}{dt} = l \dot{\vartheta}$$

$$U = m g h = m g l (1 - \cos\vartheta)$$

$$\mathcal{L} = \frac{1}{2} m l^2 \dot{\vartheta}^2 - m g l (1 - \cos\vartheta) = \mathcal{L}(\vartheta, \dot{\vartheta})$$

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\vartheta}} - \frac{\partial \mathcal{L}}{\partial \vartheta} = 0$$

$$\frac{\partial \mathcal{L}}{\partial \vartheta} = -m g l \sin\vartheta$$
$$\frac{\partial \mathcal{L}}{\partial \dot{\vartheta}} = m l^2 \dot{\vartheta}$$
$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\vartheta}} = m l^2 \ddot{\vartheta}$$



$$m l \ddot{\vartheta} + m g \sin \vartheta = 0 ; \quad \ddot{\vartheta} = -\frac{g}{l} \sin \vartheta$$

$\vartheta \ll 1$

$$\ddot{\vartheta} \approx -\frac{g}{l} \vartheta \quad \vartheta(t) = \vartheta_0 \cos(\omega t) \quad \omega = \sqrt{\frac{g}{l}}$$



Initial conditions:  $\begin{cases} \vartheta(0) = \vartheta_0 \\ \dot{\vartheta}(0) = 0 \end{cases}$

$$T = \frac{2\pi}{\omega} = 2\pi \sqrt{\frac{l}{g}}$$

$$l = l(t)$$

$$\ddot{\vartheta} \approx -\frac{g}{l(t)} \sin \vartheta$$

$$\vartheta(t) = \text{Re} \left[ \vartheta_0 e^{i\omega t} \right] = \text{Re} \left[ \vartheta_0 e^{i \int_0^t \omega dt'} \right]$$

$$\Rightarrow \omega(t) = \sqrt{\frac{g}{l(t)}}$$

Guess

$$z(t) = \text{Re} \left[ \cancel{\vartheta_0(t)} e^{i \int_0^t \omega(t') dt'} \right]$$

$$\omega(t) = \sqrt{g/l(t)}$$

$$\ddot{\vartheta} = -\frac{g}{l} \vartheta$$

$$\frac{d\vartheta}{dt} = \text{Re} \left[ \dot{\vartheta}_0 e^{i \int_0^t \omega(t') dt'} + \cancel{\vartheta_0} e^{i \int_0^t \omega(t') dt'} \cancel{i\omega(t)} \right]$$

$$\frac{d^2\vartheta}{dt^2} = \text{Re} \left[ \ddot{\vartheta}_0 e^{[ ]} + \dot{\vartheta}_0 e^{[ ]} \cdot i\omega + \vartheta_0 e^{[ ]} i\omega + \vartheta_0 e^{[ ]} (i\omega)^2 + \vartheta_0 e^{[ ]} i \frac{d\omega}{dt} \right] = \text{Re} \left[ e^{[ ]} \left( \ddot{\vartheta}_0 + 2i\omega\dot{\vartheta}_0 - \omega^2\vartheta_0 + i\vartheta_0 \frac{d\omega}{dt} \right) \right]$$

$$\operatorname{Re} \left[ e^{[ \dots ]} \left( \underbrace{\ddot{\vartheta}_0 + 2i\dot{\vartheta}_0 \omega + i\omega \dot{\vartheta}_0 - \omega^2 \vartheta_0}_{\approx 0} \right) \right] = -\omega^2 \vartheta_0$$

$$\vartheta_0 = \vartheta_0(t)$$

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$$2i\dot{\vartheta}_0 \omega = -i\omega \dot{\vartheta}_0$$

$$\omega_0 = \sqrt{\frac{p}{l_0}}$$

$$\omega(t) = \sqrt{\frac{p}{l(t)}}$$

$$\int_{\vartheta_0(0)}^{\vartheta_0(t)} \frac{d\vartheta_0}{\vartheta_0} = - \int_{\omega_0}^{\omega(t)} \frac{d\omega}{\omega}$$

$$= \ln \left[ \sqrt{\frac{\omega_0}{\omega(t)}} \right]$$

$$\ln \left[ \frac{\vartheta_0(t)}{\vartheta_0(0)} \right] = \frac{-1}{2} \ln \left[ \frac{\omega(t)}{\omega_0} \right]$$

$$\boxed{\mathcal{F}_0(t) = \mathcal{F}_0(0) \sqrt{\frac{\omega_0}{\omega(t)}}$$

$$\mathcal{F}(t) = \operatorname{Re} \left[ \mathcal{F}_0(0) \sqrt{\frac{\omega_0}{\omega(t)}} e^{i \int_0^t \omega(t') dt'} \right] =$$

$$\mathcal{F}(t) = \mathcal{F}_0(0) \sqrt{\frac{\omega_0}{\omega(t)}} \cos \left( \int_0^t \omega(t') dt' \right)$$