

# Chapter 5 - Introduction to thermonuclear fusion

## 1 Tunnel effect and fusion cross section

The effective potential seen by two particles 1 and 2 for a fusion reaction in the centre of mass frame can be represented as in figure 1. The two particles have the charge  $+Z_1e$  and  $+Z_2e$ , respectively, a reduced mass  $\mu$  and a relative kinetic energy  $\epsilon$ . The nucleus has a radius  $r_n = r_{n1} + r_{n2}$  and  $r_{tp}$  indicates the classical turning point. The potential well of the nucleus has a depth  $-U_0$  and the Coulomb barrier has a maximum height of  $V_b$ .

Classically, a particle with relative kinetic energy  $\epsilon$  cannot overcome the Coulomb barrier if  $\epsilon < V_b$ . Quantum mechanics however predicts that the particle can overcome the barrier even when  $\epsilon < V_b$  due to tunneling and can thus undergo a fusion reaction also in this case.

The solution of the Schrödinger equation for this problem shows that the effective potential seen by the wave component that carries the angular momentum  $\hbar\sqrt{l(l+1)}$  is

$$V_{eff,l}(r) = V(r) + \frac{\hbar^2 l(l+1)}{2\mu r^2} \quad (1)$$

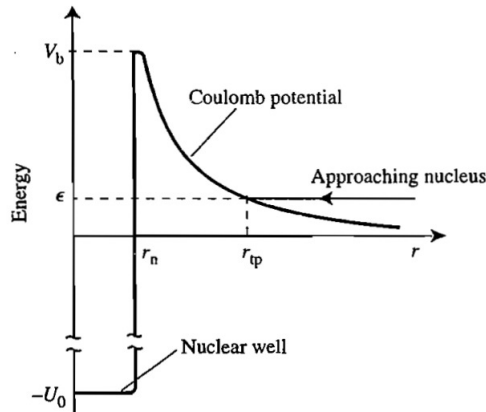


Figure 1: Effective potential in the centre of mass frame seen by a charged particle with relative kinetic energy  $\epsilon$  undergoing a fusion reaction.

where  $V(r)$  is the potential of figure 1. The angular momentum components of the wave that have  $l \neq 0$  experience a progressively increased potential barrier due to the centrifugal term  $\frac{\hbar^2 l(l+1)}{2\mu r^2}$  and have a smaller probability to undergo a fusion reaction. We can thus assume that, to a first extent, only the  $l = 0$  component contributes to the reaction, so that  $V_{eff,l}(r) = V(r)$ ; in particular, when  $r > r_n$ ,  $V(r)$  is the repulsive Coulomb potential between two charges  $+Z_1e$  and  $+Z_2e$ .

Under these conditions, the WKB theory of quantum mechanics predicts that the tunneling probability  $P$  is

$$P \approx \exp\left(-\frac{2}{\hbar} \int_{r_n}^{r_{tp}} \sqrt{2\mu(V(r) - \epsilon)} dr\right) \quad (2)$$

Evaluate  $P$  and show that, if  $r_{tp} \gg r_n$

$$P \approx \exp\left(-\sqrt{\frac{\epsilon_G}{\epsilon}}\right) \quad (3)$$

$\epsilon_G = (\pi\alpha_f Z_1 Z_2)^2 2mc^2$  is the ‘‘Gamow energy’’ and  $\alpha_f = \frac{e^2}{4\pi\epsilon_0\hbar c}$  is the fine structure constant <sup>1</sup>.

## 2 Some estimates for fusion power plants

A fusion power plant operates at the temperature  $T=15$  keV, when the  $t(d, n)\alpha$  reactivity is  $\approx 2.7 \cdot 10^{-22} \text{ m}^3/\text{s}$  so to produce 1 GW of thermal fusion power. If the plasma has an electron density  $n = 10^{20} \text{ m}^{-3}$

- How many reactions per second need to occur?
- What is the plasma volume that allows obtaining this amount of reactions per second?
- Considering that deuterium has an isotopic abundance of 0.015%, how many litres of water have to be used to produce all the deuterium in the power plant?
- What is the fraction per second of deuterium that is burnt due to the fusion reactions?

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<sup>1</sup>The following recursive formula can be useful when evaluating the integral for  $P$ . If  $J_n = \int \frac{dt}{(1+t^2)^n}$ , then

$$J_{n+1} = \frac{1}{2n} \frac{t}{(1+t^2)^n} + \frac{2n-1}{2n} J_n \quad (4)$$