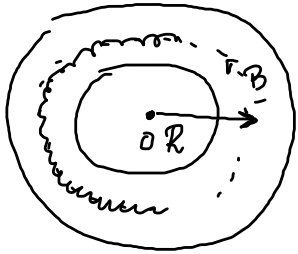


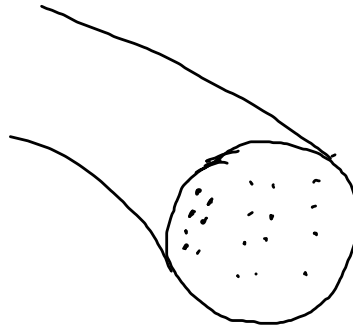
$$\underline{F} = -\mu \underline{\nabla} |\underline{B}|$$

$$\mu = \frac{m v_{\perp}^2}{2B}$$

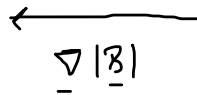
$$\underline{v} = \frac{\underline{F} \times \underline{B}}{qB^2}$$



$$B \propto \frac{1}{R}$$



O x



$$\underline{v} = -\mu \underline{\nabla} |\underline{B}| \times \underline{B}$$

Electroni qB^2

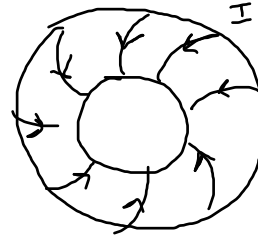
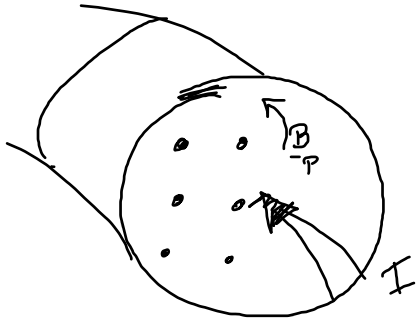
$$\underline{v} = \frac{\underline{F} \times \underline{B}}{B^2}$$

Joni $\uparrow \underline{E}$ $\frac{\underline{J} \times \underline{B}}{-\underline{E} \times \underline{B}}$

$$\frac{\underline{F} \times \underline{B}}{B^2}$$

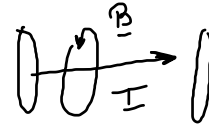
Tracce: fare avvolgere la linea di B su se stessa, durante un giro toroidale

\Rightarrow aggiungere una componente poloidale al B



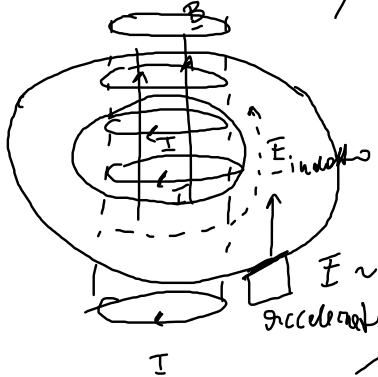
Tokamak \rightarrow
Stellarator

fare scorrere una corrente toroidale



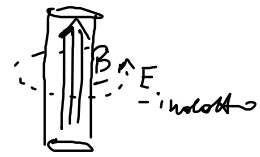
macchina intrinsecamente
stazionaria, us. complessa

Corrente toroidale \Rightarrow campo elettrico toroidale indotto



f.e.m. = $-\frac{d\phi(B)}{dt}$ legge Faraday-N-L

$\nabla \times \underline{E} = -\frac{\partial \underline{B}}{\partial t}$



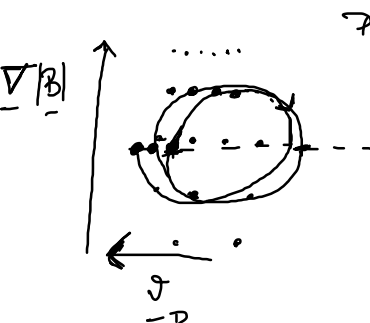
costante nel tempo $\Rightarrow \phi(B) \propto t$

$\underline{E}_{indotto}$ (f.e.m. = = =)

$\Rightarrow B \propto t$
solenoide

$\Rightarrow I_{sol} \propto t$

Dispositivo impulsato



Primo $\perp B$

$$r_L = \frac{mv_{\perp}}{qB} \propto \frac{1}{B}$$

$(-\underline{v}/|B| \times \underline{B})$ vs sinistra

$$\mathcal{L}(q_j, \dot{q}_j, t)$$

$$j = 1 \dots N$$

q_j : coordinate libere

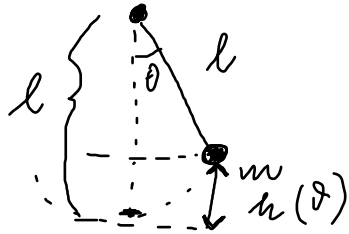
\dot{q}_j : velocità generalizzate

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{q}_j} - \frac{\partial \mathcal{L}}{\partial q_j} = 0$$

$$p_j = \frac{\partial \mathcal{L}}{\partial \dot{q}_j} \quad \text{momento canonico}$$

Se \mathcal{L} non dep. da una particolare q_k
 $\Rightarrow \frac{\partial \mathcal{L}}{\partial q_k} = 0 \Rightarrow p_k$ è costante esatta

Pendolo semplice



$$\mathcal{L} = K - U$$

eu. cinetica

eu. potenziale

$$q = \theta$$

$$K = \frac{1}{2} m v^2 = \frac{1}{2} m l^2 \dot{\theta}^2$$

$$v = l \dot{\theta}$$

$$h(\theta) = l - l \cos \theta = l(1 - \cos \theta)$$

$$U = m g h(\theta)$$

$$= m g l (1 - \cos \theta)$$

$$\mathcal{L} = \frac{1}{2} m l^2 \dot{\theta}^2 - m g l (1 - \cos \theta)$$

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\theta}} - \frac{\partial \mathcal{L}}{\partial \theta} = 0$$

$$m l^2 \ddot{\theta} + m g l \sin \theta = 0$$

$$\frac{\partial \mathcal{L}}{\partial \theta} = - m g l \sin \theta$$

$$\frac{\partial \mathcal{L}}{\partial \dot{\theta}} = m l^2 \dot{\theta}$$

$$\ddot{\vartheta} = -\frac{g}{l} \sin \vartheta$$

$$\vartheta \ll 1$$

$$\ddot{\vartheta} = -\frac{g}{l} \vartheta$$

$$\vartheta(t) = \vartheta_0 \cos(\omega t)$$

$$\begin{aligned} \vartheta(t) &= \operatorname{Re} \left[\vartheta_0 e^{i\omega t} \right] \\ &= \operatorname{Re} \left[\vartheta_0 e^{i \int_0^t \omega dt'} \right] \end{aligned}$$

$$\omega = \sqrt{\frac{g}{l}}$$

$l = l(t)$ e $l \approx \text{cost}$ su un periodo

$$\omega(t) = \sqrt{\frac{g}{l(t)}}$$

$$\ddot{\vartheta} = -\frac{g}{l(t)} \vartheta$$

Un aspetto

$$\vartheta(t) = \operatorname{Re} \left[\vartheta_0(t) e^{i \int_0^t \omega(t') dt'} \right]$$

$\xrightarrow{\sqrt{\frac{g}{l(t')}}}$