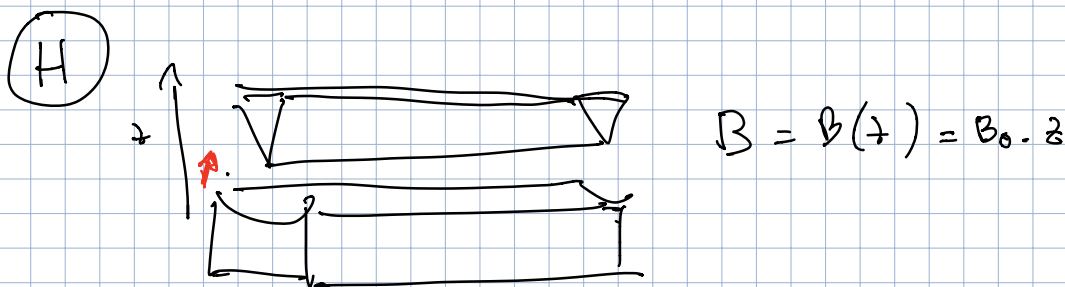


## Gas di elettroni liberi

$$E_{n_x, n_y, n_z} = C(n_x^2 + n_y^2 + n_z^2); \quad C = \frac{\hbar^2}{2m_e} \frac{\pi^2}{L^2}$$

$$\{n_x, n_y, n_z, m_s\}; \quad m_s = \pm \frac{1}{2} \quad \left\{ \begin{array}{l} \hat{S}_z \rightarrow \hbar m_s; \quad m_s = \pm \frac{1}{2} \\ \hat{S}^2 \rightarrow \hbar^2 s(s+1) \quad \left( s = \frac{1}{2} \right) \end{array} \right.$$



$$\vec{\mu} \quad E = -\vec{\mu} \cdot \vec{B} = -\mu_z \cdot B_0 \cdot z$$

$$\hat{\vec{\mu}} = \left( \frac{-e \hbar}{2m_e} \hat{L} \right) \quad E = + \frac{e \hbar}{2m_e} L_z \cdot B_0 \cdot z$$

$$F_z = - \frac{\partial E}{\partial z} = \frac{e \hbar}{2m_e} L_z \cdot B_0$$

Atomi di H nel loro stato fondamentale  $\psi_{100} \quad \left\{ \begin{array}{l} n=1 \\ l=0 \\ m=0 \end{array} \right.$

$$F_z = \frac{e}{2m_e} (\hat{L}_z) \cdot B_0$$

↳  $\hbar m = 0$



$F_z \sim L_z$  con 2 possibili autovalori  $\pm$

$$[\hat{L}_x, \hat{L}_y] = i\hbar \hat{L}_z$$

$$L^2 \rightarrow \hbar^2 \ell(\ell+1)$$

$$[\hat{L}^2, \hat{L}_z] = 0$$

$$L_z \rightarrow \hbar m_\ell ; m_\ell = -\overset{2\ell+1}{\underbrace{\ell}_{\text{valori}}}, \dots, \ell$$

$$Y_\ell^m(\theta, \varphi)$$

$$\hat{S}_x, \hat{S}_y, \hat{S}_z ; [\hat{S}_x, \hat{S}_y] = i\hbar \hat{S}_z$$

$$S^2 \rightarrow \hbar^2 s(s+1)$$

$$S_z \rightarrow \hbar m_s ; m_s = -\overset{1}{\underbrace{s}_{\text{valori}}}, \dots, s$$

$$[\hat{S}^2, \hat{S}_z] \rightarrow \text{aut. comuni } \chi_{m_s}^s$$

$$S^2 \chi_{m_s}^s = \hbar^2 s(s+1) \chi_{m_s}^s \quad s = \frac{1}{2}$$

$$L_z \chi_{m_s}^s = \hbar m_s \chi_{m_s}^s$$



$$\sum_{n_x, n_y, n_z, n_s} \epsilon_n \rightarrow \epsilon_\alpha$$

$$\left( n_{n_x, n_y, n_z, n_s} \right)_{FD} = \frac{1}{e^{\beta(\epsilon_{n_x, n_y, n_z, n_s} - \mu)} + 1}$$

$$P_{n_x, n_y, n_z, n_s} = \frac{e^{-\beta \epsilon_{n_x, n_y, n_z, n_s}}}{Z_T}$$

$$\epsilon_{n_x, n_y, n_z, n_s} = C \cdot (n_x^2 + n_y^2 + n_z^2) \rightarrow \mu^2$$

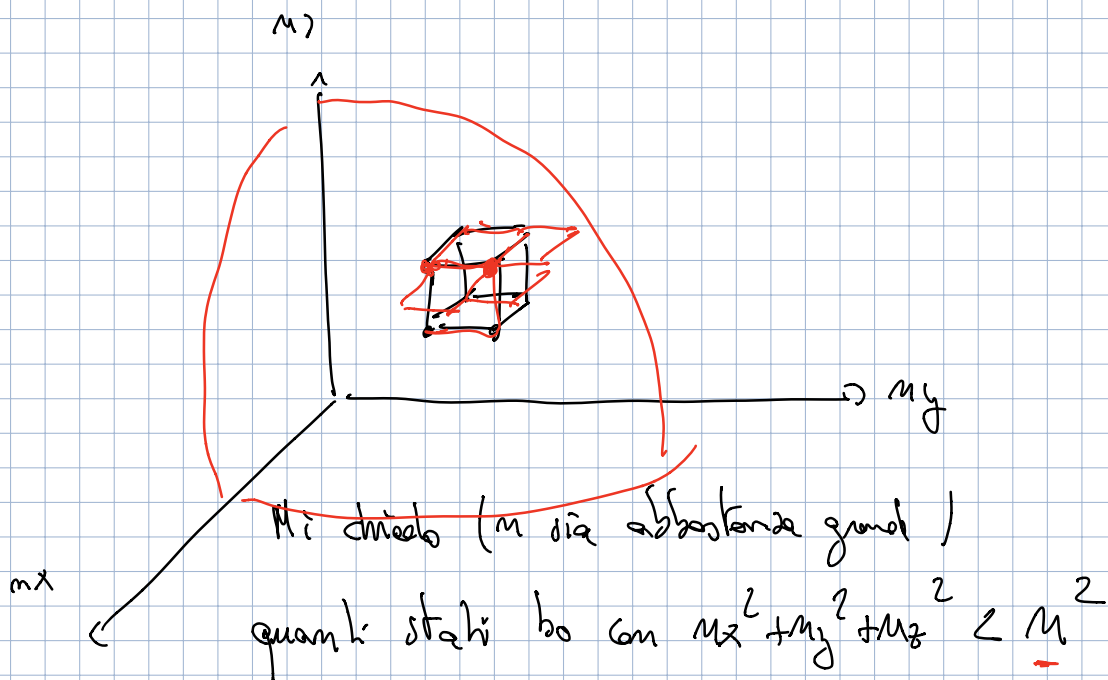
$\Sigma$  für  $H$

$$P_{n_x, n_y, n_z, n_s} = \frac{e^{-\beta \epsilon_{n_x, n_y, n_z, n_s}}}{\sum_{n_x, n_y, n_z, n_s} e^{-\beta \epsilon_{n_x, n_y, n_z, n_s}}}$$

$$\Sigma = -\frac{R \ln C}{n^2}$$

$$P_n = \frac{e^{-\beta \epsilon_n} \cdot g(n)}{\sum_n e^{-\beta \epsilon_n} g(n)} \rightarrow 2n^2$$





$$\frac{V}{m} = \frac{\frac{4}{3} \pi m^3}{\frac{4}{3} \pi m^3} / V_{\text{cubo}} = 1 \quad \equiv \text{Numero di punti} \\
 \text{con } m_x^2 + m_y^2 + m_z^2 < m^2$$

Quanti stati ho nell'intervallo di energia  $(m, m+dm)$

$$V_{m+dm} = \frac{4}{6} \pi (m+dm)^3 \\
 \hookrightarrow m^3 + 3m^2 dm + \text{termini in } dm^2 \text{ o } dm^3$$

$N(m) \cdot dm =$  numero di punti tali per cui  $(m_x, m_y, m_z)$  sono  
 tali che  $m_x^2 + m_y^2 + m_z^2 \in U_{dm}(m) = V_{m+dm} - V_m$

$$N(dm) = \frac{3\pi m^2 dm}{6} = \frac{\pi m^2 dm}{2} \quad \varepsilon = cm^2$$

$$m^2 = \frac{\varepsilon}{c}$$

$$\frac{\pi m^2 dm}{2} = \frac{\pi \varepsilon}{2c} \cdot \frac{1}{2\sqrt{\varepsilon}} \cdot \frac{1}{\sqrt{c}} d\varepsilon$$

$$m = \sqrt{\frac{\varepsilon}{c}}$$

$$= \frac{\pi \sqrt{\varepsilon}}{4 c^{3/2}} d\varepsilon$$

$$dm = \frac{1}{2\sqrt{\varepsilon} \cdot \sqrt{c}} d\varepsilon$$

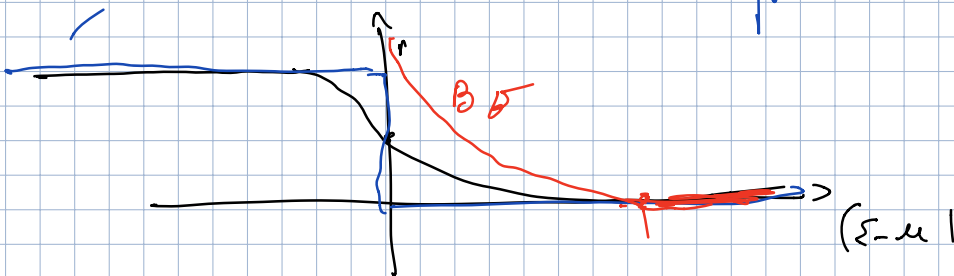
↳ numero di stati con energia compresa tra  $\varepsilon$  e  $\varepsilon + d\varepsilon$

Anzi, se considero lo spin  $g(\varepsilon) = \frac{\pi \sqrt{\varepsilon}}{2 c^{3/2}} d\varepsilon$

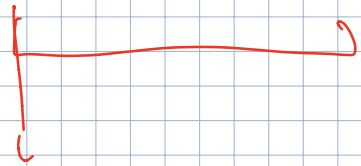
$$\left( \int_{m_x, m_y, m_z, m_s} f \right)_{FD}$$

$$\frac{1}{e^{\beta(\varepsilon_{m_x, m_y, m_z, m_s} - \mu)} + 1}$$

$$\frac{1}{e^{\beta(\varepsilon - \mu)} - 1}$$



$$\sum_{n_x, n_y, n_z} \left( f_{n_x, n_y, n_z} \right)_{FD} = \langle N \rangle = N$$



$$2 \sum_n g(n) \left( f_n \right)_{FD}$$

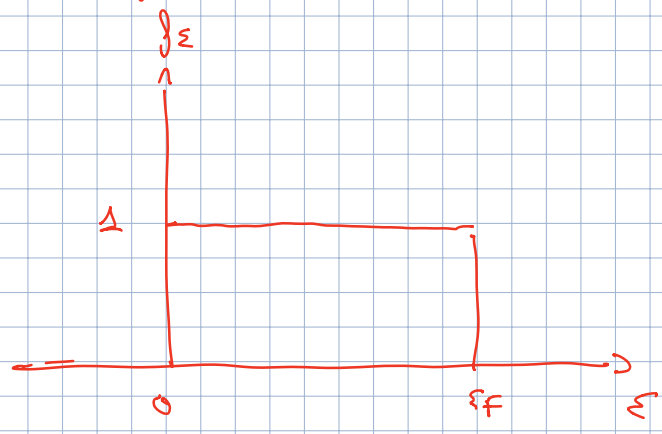
$\frac{\pi}{2} \frac{\sqrt{\epsilon}}{c^{3/2}} \quad c = \frac{h^2}{2m} \left( \frac{\pi}{L} \right)^2$

$$\int_0^{+\infty} g(\epsilon) \cdot f_{FD}(\epsilon) d\epsilon = \langle N \rangle$$

$$\frac{1}{e^{\beta(\epsilon - \mu)} + 1} ; \mu = \mu(V, N, T)$$

Possiamo immaginare  $V$  fissato,  $N$  fissato  $\Rightarrow \mu = \mu(T)$ .

Definiamo  $\epsilon_F = \mu(0)$





$$N = \int_G^{\infty} g(\epsilon) f_{FD}(\epsilon) d\epsilon \stackrel{T \rightarrow 0}{\approx} \int_0^{\epsilon_F} g(\epsilon) d\epsilon$$

$$N = \int_0^{\epsilon_F} \frac{\pi}{2} \frac{\sqrt{\epsilon}}{C^{\frac{3}{2}}} d\epsilon = \frac{\pi}{2} \frac{1}{C^{\frac{3}{2}}} \frac{2}{3} \epsilon_F^{\frac{3}{2}}$$

$$N = \frac{\pi}{3 \cdot C^{\frac{3}{2}}} \epsilon_F^{\frac{3}{2}}$$

$$C = \frac{\hbar^2}{2m} \frac{\pi^2}{C^2}$$

$$C^{\frac{3}{2}} = \frac{\hbar^3 \pi^3}{(2m)^{\frac{3}{2}} \sqrt{V}}$$

$$N = \frac{\pi}{3 \hbar^3 \pi^{\frac{3}{2}}} (2m)^{\frac{3}{2}} \cdot V \epsilon_F^{\frac{3}{2}}$$

$$\rho = \frac{N}{V}$$

$$\epsilon_F^{\frac{3}{2}} = \frac{3 \hbar^3 \pi^{\frac{3}{2}} \rho}{(2m)^{\frac{3}{2}}}; \quad \epsilon_F = \underbrace{\left( \frac{3 \pi^2 \rho}{k_F^3} \right)^{\frac{2}{3}}}_{k_F^2} \cdot \frac{\hbar^2}{2m_e}$$

$\rho$  el. and metallo  $\approx 10^{23} / \text{cm}^3 \rightarrow$  typische  $\epsilon_F \approx 5 \text{ eV}$



$\epsilon_F \approx 5 \text{ eV}$       $k_B T = 5 \text{ eV}$       $\text{in } T = 60000 \text{ K}$

A temperature ragionevoli:  $\mu \approx \epsilon_F$  (non esattamente, ma  $\int_{FD} \approx$  gradino sempre connessi).

Energia media a  $T \approx 0$

$$\langle U(0) \rangle = \int_0^{\epsilon_F} g(\epsilon) \cdot \epsilon \cdot d\epsilon$$

$$= \frac{\pi}{2} \frac{1}{C^{\frac{3}{2}}} \int_0^{\epsilon_F} \epsilon^{\frac{3}{2}} = \frac{\pi}{2} \frac{1}{C^{\frac{3}{2}}} \cdot \frac{2}{5} \epsilon_F^{\frac{5}{2}}$$

$$= \frac{\pi}{2} \frac{1}{C^{\frac{3}{2}}} \frac{1}{5} \epsilon_F \cdot \epsilon_F^{\frac{3}{2}} = \frac{3}{5} N \epsilon_F$$

$$N = \frac{\pi}{3 \cdot C^{\frac{3}{2}}} \epsilon_F^{\frac{3}{2}}$$

$$\Rightarrow \epsilon_F^{\frac{3}{2}} = \frac{3 N C^{\frac{3}{2}}}{\pi}$$

$$\Rightarrow \frac{U(0)}{N} = \frac{3}{5} \epsilon_F$$