

Chapter 1 - Introduction to plasma physics

1 Rutherford scattering

Consider a Rutherford scattering process between two particles having velocities $\mathbf{v}_1, \mathbf{v}_2$, charges q_1, q_2 and masses m_1, m_2 , respectively.

- a) Show that, in the centre of mass frame, the equation of motion is given by

$$\mu \ddot{\mathbf{r}} = \frac{q_1 q_2}{4\pi\epsilon_0 r^2} \hat{\mathbf{r}} \quad (1)$$

where \mathbf{r} is the relative position vector.

- b) Define the angular momentum $\mathbf{L} = \mu \mathbf{r} \times \dot{\mathbf{r}}$. By taking its derivative with respect to the time, show that this is a constant vector. In particular, note that this implies that the scattering process must occur in the plane where \mathbf{r} and $\dot{\mathbf{r}}$ lie.
- c) Using the geometry shown in 1, show that $|\mathbf{L}| = \mu b v_\infty = \mu r^2 \dot{\phi}$, from which $\dot{\phi} = b v_\infty / r^2$.
- d) Call \mathbf{v}_i and \mathbf{v}_f the initial and final velocity of the charged particle shown in the figure and note that $|\mathbf{v}_i| = |\mathbf{v}_f|$. From symmetry, the x component of the velocity before and after scattering must be unchanged, whereas the y component must flip. Starting from these considerations, evaluate the change of the y velocity component before and after the collision (Δv_y) as a function of v_{yi} .

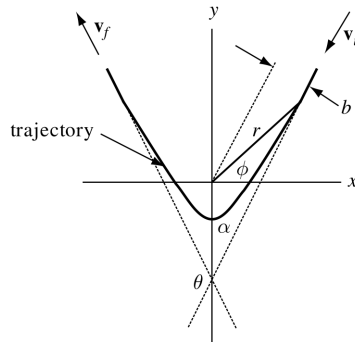


Figure 1: Scattering geometry in the centre of mass frame. The scattering centre is in the origin and θ denotes the scattering angle.

- e) Starting from the y component of equation 1, find a relation between dv_y and $d \cos \phi$. (Hint: consider the conservation of the angular momentum to write dt as a function of $d\phi$). Call ϕ_i and ϕ_f the initial and final values of ϕ . By integrating dv_y , find Δv_y for the whole collision and as a function of ϕ_i and ϕ_f .
- f) Using the figure, find the relation between the angles θ and α and between the angles ϕ_i, ϕ_f and α . By combining the different formulas for Δv_y found above, show that

$$\tan \frac{\theta}{2} = \frac{q_1 q_2}{4\pi\epsilon_0 \mu b v_\infty^2} \quad (2)$$

and find an approximation for small angle collisions.

2 Debye shielding

We aim at solving the 3-dimensional Poisson equation for the electrostatic potential $\phi(\mathbf{r})$ of a point charge q_T in a uniform plasma of ions and electrons

$$\nabla^2 \phi - \frac{1}{\lambda_D^2} \phi = -\frac{q_T}{\epsilon_0} \delta(\mathbf{r}) \quad (3)$$

where λ_D is the Debye length and δ is the Dirac-delta.

- a) Starting from the known solution of the Poisson equation

$$\nabla^2 \phi = -\frac{q_T}{\epsilon_0} \delta(\mathbf{r}) \quad (4)$$

of a charge q_T in vacuum, which is

$$\phi(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{q_T}{r} \quad (5)$$

show that we can write

$$\delta(\mathbf{r}) = -\frac{1}{4\pi} \nabla^2 \left(\frac{1}{r} \right) \quad (6)$$

- b) Using this result in (3), evaluate the laplacian operator in spherical coordinates and find a solution for (3) of the type

$$\phi(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{g(r)}{r} \quad (7)$$

In particular, work out an equation for $g(r)$ e show that the solution is $g(r) = \exp(-r/\lambda_D)$.